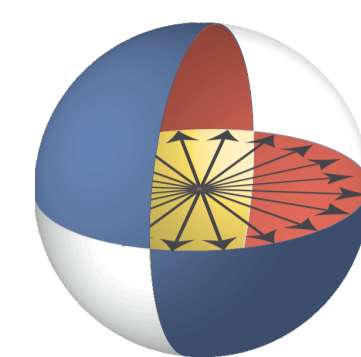


# Simulating quantum circuits by classical circuits

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## Abstract

In a recent breakthrough, Bravyi, Gosset, and König (BGK) [1] proved that shallow quantum circuits cannot be "simulated" by shallow classical circuits. Indeed, they show a lower bound of log depth. **Our work addresses the upper bound question.**

We first explicitly define their implicit notion of simulation which we call "possibilistic simulation" (see right). Essentially, we say a classical circuit simulates a quantum circuit if, over all inputs, the output of the classical circuit is a possible output of the quantum circuit.

In this sense, classically simulating the BGK quantum circuits that solve their Hidden Linear Function (HLF) problem is equivalent to classically solving HLF [1].

We then show the following two incomparable results.

**Result 1.** Any quantum circuit of depth  $D$  with Clifford gates and  $t$   $T$  gates can be simulated in depth

$$O(D + t)$$

where  $O$  conceals a constant independent of the quantum circuit, in particular, its number of qubits.

**Result 2.** The BGK quantum circuits can be simulated in complexity class  $NC^2$ , i.e. by classical circuits of **log squared** depth and polynomial size (in number of qubits).

## Discussion

Result 1 follows from the construction to the right. While Result 1 is inspired by Gottesman-Knill (GK) and an extension by Bravyi and Gosset [2], it is not directly implied by them. For example, GK does not directly imply Result 1 with  $t=0$  because a usual GK simulator would "measure"  $n$  bits sequentially using depth  $O(n)$ . In fact, one way to interpret Result 1 is that it parallelises classical simulators in the context of possibilistic simulation.

Result 2 follows because HLF can be solved by finding a matrix kernel and then solving a linear equation [1]. But these tasks are in complexity class  $NC^2$ , e.g. [3].

**Correction.** The current (to-be-updated) arXiv paper costs the depth of a "switchboard circuit" in its "Construction B" as  $O(\log t)$ . This should be  $O(t)$  as pointed out to me by Luke Schaeffer. The main consequence is that "Result 4" in the paper, saying log depth separation is the maximum achievable, is revised to be Result 1 above.

**Ongoing work for v3.** Extending Result 1 to parallelise other simulation techniques, e.g. tensor network based.

## Possibilistic simulation definition

**Definition 1.** We make the following definitions for circuits with  $n$  variable input bits and  $m$  output bits.

- A relation on Cartesian product  $\{0, 1\}^n \times \{0, 1\}^m$  is a subset  $\mathcal{R} \subset \{0, 1\}^n \times \{0, 1\}^m$ .
- A quantum circuit  $Q$  on  $n$  input qubit lines and measured on  $m$  qubit lines in the computational basis defines a relation  $\mathcal{R}(Q) \subset \{0, 1\}^n \times \{0, 1\}^m$  by:

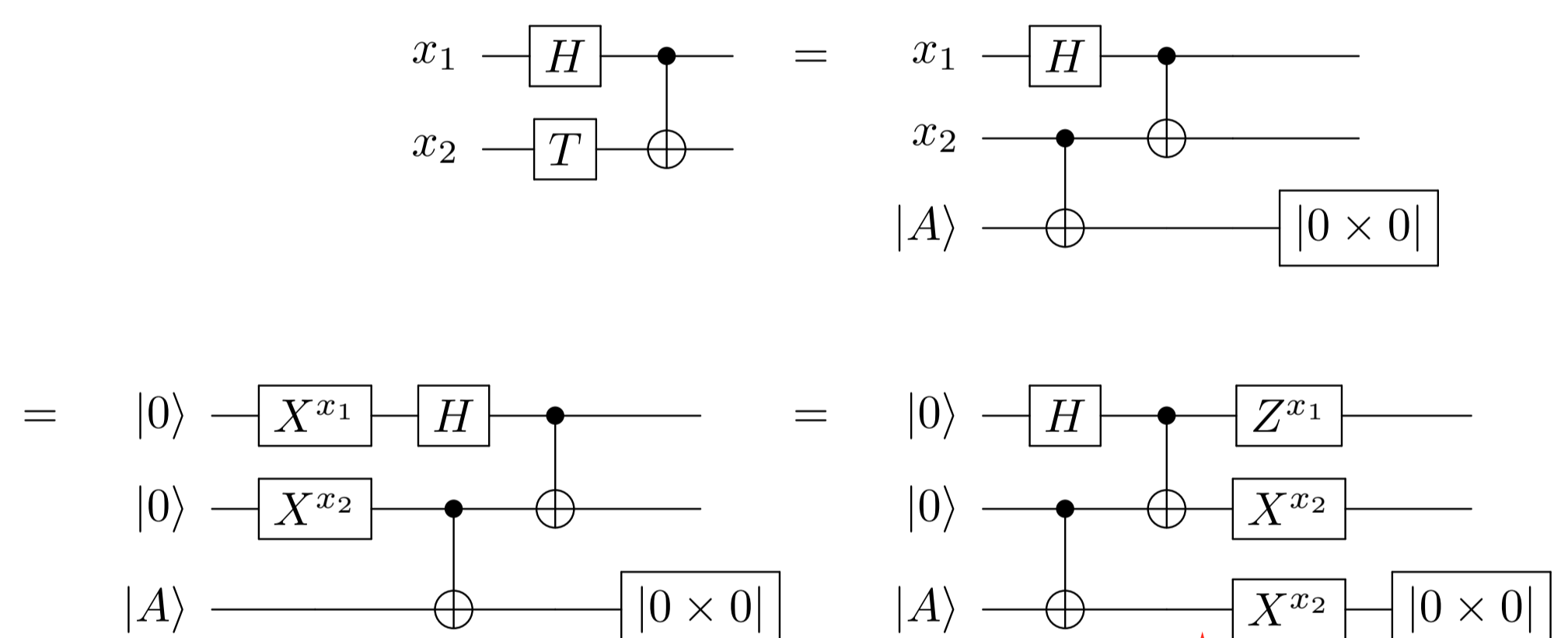
$$(x, y) \in \mathcal{R}(Q) \iff \langle y | Q | x \rangle \neq 0. \quad (1)$$

- Let  $C : \{0, 1\}^n \rightarrow \{0, 1\}^m$  be a classical circuit, and  $\mathcal{R}$  a relation on  $\{0, 1\}^n \times \{0, 1\}^m$ . We say  $C$  possibilistically simulates  $\mathcal{R}$  if:

$$(x, C(x)) \in \mathcal{R}, \text{ for all } x \in \{0, 1\}^n. \quad (2)$$

## Construction for Result 1

Idea: use circuit identities like those in the example below together with free pre-computation to construct classical Clifford+ $T$  simulator.



$$\text{with } |A\rangle = \frac{1}{\sqrt{2}}(|0\rangle + \omega|1\rangle) \text{ where } \omega = e^{i\pi/4}$$

Now we can *precompute* the 3-qubit state at the red arrow:

$$|\psi\rangle = \frac{1}{2}(|000\rangle + |110\rangle + \omega|001\rangle + \omega|111\rangle)$$

which is independent of the input bits  $x_i$ .

To simulate the circuit, we may first use  $x_2$  to determine which of the two 2-qubit "sectors" of  $\psi$  to project to. Suppose  $x_2 = 0$ , then we should project to the sector

$$|\psi_0\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

Let  $s$  be a bit-string in the support of  $\psi_0$ , say  $00$ , which can be precomputed. We can then apply the "classical version" of  $Z_1^{x_1} X_2^{x_2}$  to  $s$  to classically output a possible output of the quantum circuit. This completes our construction.

The  $t$  in Result 1 comes from a "switchboard circuit" that carries out projections onto  $2^t$  sectors in general. The  $D$  is because at most  $2^D$  bits  $x_i$  can appear in the exponent of each Pauli operator on the right-hand-side of the last circuit above. The  $2^t$  (resp.  $2^D$ ) becomes linear in  $t$  (resp.  $D$ ) using binary trees of various kinds.

## References

- [1] S. Bravyi, D. Gosset, and R. König, Science **362**, 208 (2018).
- [2] S. Bravyi and D. Gosset, Phys. Rev. Lett. **116**, 250501 (2016).
- [3] K. Mulmuley, Combinatorica **7**, 101 (1987).