

# Lecture 16

**Definition 3** (Mixed quantum state). Let  $d \in \mathbb{N}$ . A  $d$ -dimensional mixed (quantum) state is a matrix  $\rho \in \mathbb{C}^{d \times d}$  that is Hermitian, positive semi-definite (PSD), and has trace 1.<sup>3</sup>

Quantum states  $|\psi\rangle \in \mathbb{C}^d$  are often called “pure quantum states” and can be viewed as mixed quantum states of rank 1 via  $|\psi\rangle \leftrightarrow |\psi\rangle\langle\psi|$ .

**Fact 4.** In quantum information, a probability distribution over quantum states  $(p_1, |\psi_1\rangle), \dots, (p_k, |\psi_k\rangle)$  is described as the matrix  $\sum_{i=1}^k p_i |\psi_i\rangle\langle\psi_i|$ , which can be checked to be a mixed quantum state according to the definition above. The converse also holds in the sense that every mixed quantum state  $\rho$  can be written in the form  $\sum_{i=1}^k p_i |\psi_i\rangle\langle\psi_i|$ , by considering its spectral decomposition.

*Comment:* the mapping from distribution to matrix is not invertible but it does not destroy information, since the information content of  $(p_i, |\psi_i\rangle)$  is  $\sum_i p_i |\psi_i\rangle\langle\psi_i|$ . Put another way: you cannot distinguish between  $(p_i, |\psi_i\rangle)$  and  $(p'_i, |\psi'_i\rangle)$  using any procedure if their corresponding mixed states are the same. E.g.,  $(\frac{1}{2}, |0\rangle), (\frac{1}{2}, |1\rangle)$  and  $(\frac{1}{2}, |+\rangle), (\frac{1}{2}, |-\rangle)$ .

The effect of a quantum measurement on a mixed quantum state is consistent with the effect of a quantum measurement on (pure) quantum states.

**Definition 4** (Schatten  $p$ -norms). Let  $p \in [1, \infty)$  and  $d \in \mathbb{N}$ . The Schatten  $p$ -norm of  $A \in \mathbb{C}^{d \times d}$  is defined to be

$$\|A\|_p := [\text{Tr}[(A^\dagger A)^{p/2}]^{1/p} \tag{10}$$

The Schatten  $\infty$ -norm of  $A$  is defined to be the spectral norm of  $A$ , which coincides with  $\lim_{p \rightarrow \infty} \|A\|_p$ .

**Definition 5** (Fidelity between mixed quantum states). For  $\rho, \sigma \in \mathbb{C}^{d \times d}$  that are mixed quantum states, the fidelity between  $\rho$  and  $\sigma$  is defined to be

$$F(\rho, \sigma) := \|\sqrt{\rho}\sqrt{\sigma}\|_1^2, \tag{11}$$

*Comment:* 04/03/26: I think I missed a square in the definition I gave of  $F$  in class. The above is the correct one.

*Comment:* have  $F(\rho, \sigma) = \text{tr}[\sqrt{\rho}\sqrt{\sigma}]^2$ , where  $|X| := \sqrt{X^\dagger X}$ . (This definition is symmetric.) Note: for  $C$  Hermitian of the form  $\sum_i \lambda_i |v_i\rangle\langle v_i|$  with  $\lambda_i \geq 0$ ,  $\sqrt{C}$  is defined to be  $\sum_i \sqrt{\lambda_i} |v_i\rangle\langle v_i|$ . Also  $F(A, B) = F(B, A)$  follows from  $\text{tr}[|X|] = \text{tr}[|X^\dagger|]$  for any  $X \in \mathbb{C}^d$  – think SVD.

**Lemma 2** (Pretty Good Measurement). Let  $\rho_1, \dots, \rho_N \in \mathbb{C}^{d \times d}$  be mixed quantum states. Then there exists a procedure that succeeds in distinguishing them with probability at least  $1 - N \sqrt{\max_{a \neq b} F(\rho_a, \rho_b)}$ .

*Proof.* Omitted. See [Harrow and Winter’06] (which applies von Neumann’s minimax theorem to [Barnum and Knill’00]).  $\square$

**Lemma 3** (Hölder’s inequality for Schatten  $p$ -norms). Let  $p \in [1, \infty]$  and  $p^* \in [1, \infty]$  satisfy  $1/p + 1/p^* = 1$ . Let  $d \in \mathbb{N}$ . Then, for  $A, B \in \mathbb{C}^{d \times d}$ , we have

$$\|AB\|_1 \leq \|A\|_p \|B\|_{p^*}. \tag{12}$$

*Proof.* Omitted, non-trivial. Variant more often seen is  $|\langle A, B \rangle| \leq \|A\|_p \|B\|_{p^*}$ , where  $\langle A, B \rangle := \text{tr}[A^\dagger B]$ . Here’s a reduction to that variant. Let  $r := \text{rk}(A)$ . Write

$$AB = \sum_{i=1}^r \sigma_i |u_i\rangle\langle v_i|, \tag{13}$$

so  $\|AB\|_1 = \sum_{i=1}^r \sigma_i$ .

Let  $U \in \mathbb{C}^{d \times d}$  unitary be such that  $U^\dagger |v_i\rangle = |u_i\rangle$ . Then

$$AB = \sum_{i=1}^r \sigma_i U^\dagger |v_i\rangle\langle v_i|. \tag{14}$$

So

$$\text{tr}[UAB] = \sum_{i=1}^r \sigma_i. \tag{15}$$

But

$$\text{tr}[UAB] = \langle A^\dagger U^\dagger, B \rangle \leq \|A^\dagger U^\dagger\|_p \|B\|_{p^*} = \|A\|_p \|B\|_{p^*}, \tag{16}$$

where the last equality uses the unitary invariance of the Schatten  $p$ -norm and its invariance under conjugate transpose.  $\square$

<sup>3</sup>In fact, the usual PSD definition, i.e.,  $u^\dagger \rho u \geq 0$  for all  $u \in \mathbb{C}^d$ , implies  $\rho$  has to be Hermitian; so the Hermitian part of the definition is somewhat redundant.

**Remark 3.** Further ingredients: von Neumann’s trace inequality:  $|\operatorname{tr}[AB]| \leq \sum_{i=1}^d \sigma_i(A)\sigma_i(B)$  (see [Stackexchange post]) and the normal Hölder’s inequality.

The next lemma follows from Hölder’s inequality with  $p = 2$ .

**Lemma 4.** *Let  $\rho, \sigma$  be mixed quantum states. Then  $F(\rho, \sigma) \leq \operatorname{tr}[\Pi_\rho \cdot \sigma]$ , where  $\Pi_\rho$  is the orthogonal projector onto the support of  $\rho$ .*

*Proof.* Have  $F(\rho, \sigma) = \|\sqrt{\rho}\sqrt{\sigma}\|_1^2 = \|\sqrt{\rho} \cdot \Pi_\rho \sqrt{\sigma}\|_1^2 \leq \|\sqrt{\rho}\|_2^2 \cdot \|\Pi_\rho \sqrt{\sigma}\|_2^2 = \operatorname{tr}[\rho] \cdot \operatorname{tr}[\sqrt{\sigma}\Pi_\rho\sqrt{\sigma}] = \operatorname{tr}[\Pi_\rho\sigma]$ . □

**Lemma 5.** *Let  $G$  be a finite group,  $H, H' \leq G$  (subgroups of  $G$ ), and  $g, g' \in G$ . Then*

$$|gH \cap g'H'| = \begin{cases} |H \cap H'| & \text{if } g^{-1}g' \in HH', \\ 0 & \text{otherwise.} \end{cases} \tag{17}$$

*Proof.* See Homework 2. □

**Lemma 6.** *Let  $G$  be a finite group. Then  $G$  has at most  $|G|^{\log_2 |G|} = 2^{(\log_2 |G|)^2}$  subgroups.*

*Proof.* For distinct elements  $g_1, \dots, g_k$ , the subgroup they generate is the subgroup formed by all words of the form  $g_{i_1}^{b_1} \dots g_{i_l}^{b_l}$ , where  $b_j \in \{-1, 1\}$  and  $i_j \in [k]$  (the empty word corresponds to  $e$ ). Written  $\langle g_1, \dots, g_k \rangle$ . We say the  $g_i$ s are independent generators of  $\langle g_1, \dots, g_k \rangle$  if  $g_i \notin \langle \{g_j : j \neq i\} \rangle$  for all  $i$ .

Every subgroup can be specified by a finite number  $k$  of independent generators. Moreover, the number of elements in the subgroup is at least  $2^k$ . This can be seen by induction: true for  $k = 0$  (the empty set generates the subgroup  $\{e\}$ ), suppose true for  $k - 1$ , then group axioms imply that  $\{g_k h \mid h \in \langle g_1, \dots, g_{k-1} \rangle\}$  has at least  $2^{k-1}$  elements, none of which are in  $\langle g_1, \dots, g_{k-1} \rangle$  as  $g_k \notin \langle g_1, \dots, g_{k-1} \rangle$ . So  $k \leq \log_2 |G|$ .

Since  $e$  cannot be an independent generator, the number of choices of independent generators is at most

$$\binom{|G| - 1}{\leq \log_2 |G|} \leq (|G| - 1 + 1)^{\log_2 |G|}, \tag{18}$$

where the inequality follows by considering the number of ways of choosing one or zero elements from  $|G| - 1$  elements.<sup>4</sup> □

All of the above serves as preliminaries for the following result.

**Proposition 1.** *Let  $G$  be a finite group. Let  $\mathcal{H}$  be the set of all subgroups of  $G$ . Then  $Q(\operatorname{HSP}(G, \mathcal{H})) = O(\log^2(|G|))$ .*

*Proof.* We will cover this following Chapter 10.3 of [AMC]. □

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<sup>4</sup>In fact, the inequality is loose and also holds if the LHS counts ordered tuples rather than subsets. However, it does yield the right “order-of-magnitude”, see Corollary 1.6 in this paper and the surrounding text.